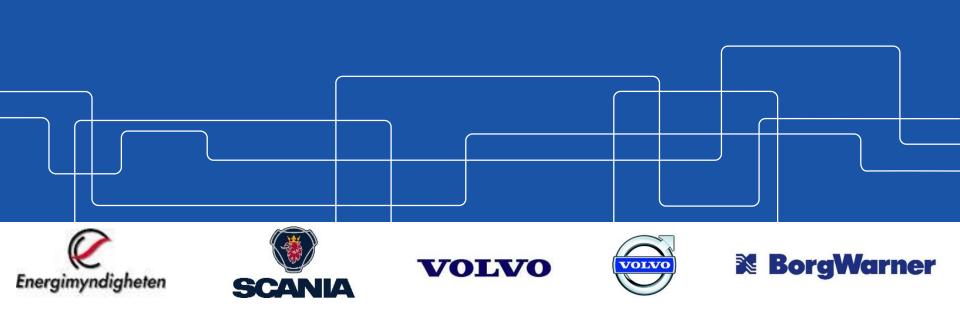


On the Aerodynamically Generated Noise in Turbocharger Compressors

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Presentation overview



- Project overview
- Research questions and approach
- Computational setup
- Results
- Summary and conclusions
- Future plans

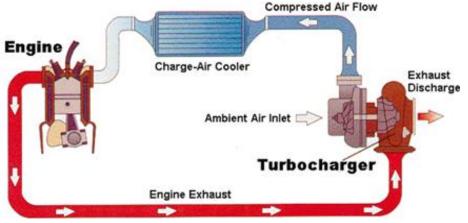


Project overview



□ ICE downsizing in automotive industry

- Regulations* and fuel economy
- Compressor noise



http://www.marine-knowledge.com/wp-content/uploads/2013/10/Turbocharger-Working.png

Project's objective: physics-based understanding of the aerodynamically generated noise in turbocharger compressors

- Investigation: CFD/CAA: RANS, URANS, LES
- Practical relevance: Propagation in real systems: installation effects



Research questions & approach



Research questions:

- Which are the dominant acoustic sources related to turbocharger noise?
- What are the implications in sound propagation and resonance of installations of turbochargers?
- How to mitigate noise produced by automotive turbochargers?
- What are the limitations of the RANS-based formulations within the context? Trends, confidence intervals?

Approach:

- RANS: compressor performance and acoustic models (advantages/limitations)
- LES (at a later stage): physics-based understanding



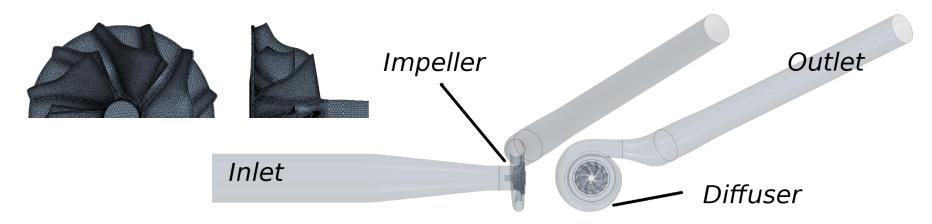
Computational setup



Compressor system

- Governing equations:
- Turbulence modelling:
- Solver:
- Discretisation:
- Mesh:

- Continuity, Momentum, Energy, Equation of State
 - SST k-ω Coupled Flow (density based)
- 2nd order upwind
- Polyhedral, ~4.5 mln cells, circumferential time averaged interface, moving reference frame

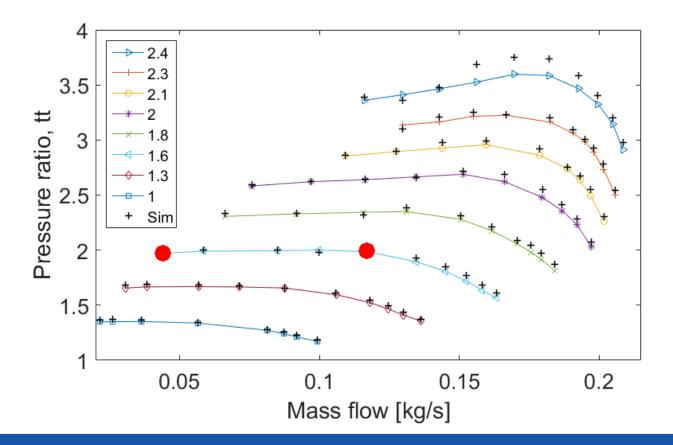




Results: Compressor performance map



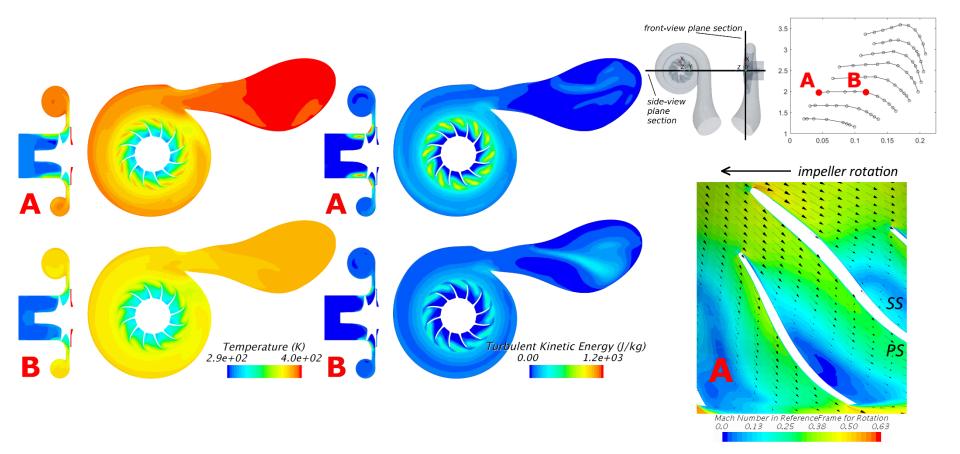
Pressure Ratio: RANS vs. Gas-stand data





Results: Temperature and TKE





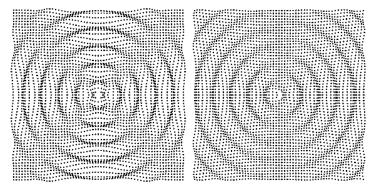


Broadband noise source models



Proudman model, Curle model (Lighthill's analogy)

- RANS turbulence modelling, acoustic source information
- Turbulence-generated flow noise
- Proudman model: quadrupole sources
- Curle model: dipole sources

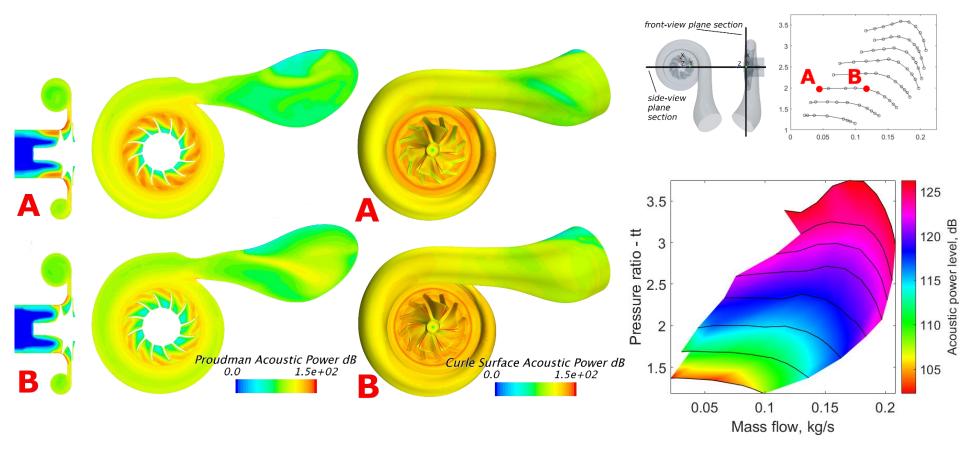


https://www.acs.psu.edu/drussell/



Results: Proudman and Curle models







Summary and conclusions



Steady-state based investigation of compressor noise

- Specific compressor performance parameters validated to experimental data
- Acoustic models in STAR-CCM+ utilised to localise areas of noise generation and estimate acoustic power level dB

Conclusions

- Suitable computational setup (compressor map)
- Noise generation areas: backflow/separation areas, volute tongue, leading edges
- Noise map: generated acoustic power proportional rpm, higher on surge/choke lines

Applicability of broadband noise source models to predict acoustic behaviour?



Future plans



Development towards correlation of sources and noise

- Refine grids for unsteady simulations (URANS → LES) on selected operating points
- Assess the validity of the RANS acoustic predictions
- Installation effects: upstream/downstream
- Compressible LES: quantification of acoustic sources and correlation of acoustic sources with far-field noise → Noise reduction technologies



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competence Center for Gas Exchange

"Charging for the future"











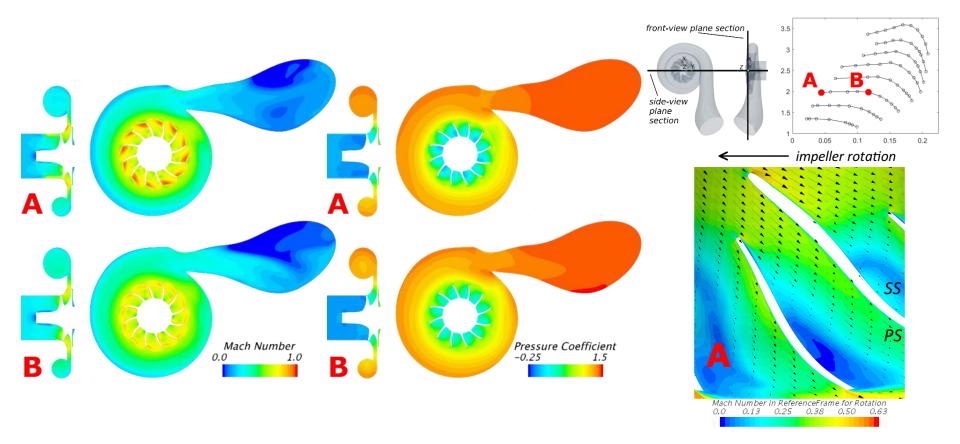


Appendix



Results: Mach number and Cp







Appendix: Curle model

The Curle surface integral is:

$$\rho'(\overline{x},t) = \frac{1}{\left(4\pi a_0^3\right)} \int_S \left[\frac{(\overline{x}-\overline{y})}{r^2} \frac{\partial p}{\partial t} \left(\overline{y}, t - \frac{r}{a_0}\right)\right] \overline{n} \, \mathrm{d}\, S(\overline{y}) \tag{954}$$

where:

- $t r / a_0$ is the emission time
- *p* the surface pressure
- ρ' the acoustic pressure
- a_0 the far-field sound speed.

On the assumption of small perturbations and an adiabatic problem, then:

$$\frac{p}{\rho^{\gamma}} = ct$$
(955)

which can be used to relate variations in acoustic pressure with density perturbations:

$$\mathbf{p'} = a_0^2 \mathbf{\rho'} \tag{956}$$



Appendix: Curle model

Then Eqn. (954) becomes:

$$\mathbf{p}'(\overline{x},t) = \frac{1}{(4\pi a_0)} \int_{S} \left[\frac{(\overline{x}-\overline{y})}{r^2} \frac{\partial p}{\partial t} \left(\overline{y}, t - \frac{r}{a_0} \right) \right] \overline{n} \, \mathrm{d}\, S(\overline{y})$$
(957)

The acoustical directional intensity per unit surface of the solid body on the far field prediction is approximated with:

$$\overline{\mathbf{p}'^2} \approx \frac{1}{16\pi^2 a_0^2} \int_{S} \frac{(\cos\theta)^2}{r^2} \left[\frac{\partial p}{\partial t} \left(\overline{y}, t - \frac{r}{a_0} \right) \right]^2 A_c(\overline{y}) \,\mathrm{d}\,S(\overline{y})$$
(958)

where:

 A_c is the correlation area

$$r = (\overline{x} - \overline{y})$$

heta is the angle between r and the $\overline{ec{n}}$ wall-normal direction.

The measure of the local contribution to acoustic power per unit surface area can be computed from:

$$SAP = \frac{1}{\rho_0 a_0} \left[\int_0^{(2\pi)} \int_0^{\pi} \overline{\mathbf{p}'}^2 r^2 \sin\theta \ d\theta \ d\gamma \right] = \int_S I(\overline{\mathbf{y}}) \, \mathrm{d}\, S(\overline{\mathbf{y}}) \qquad \qquad = \int_S \frac{A_c(\overline{\mathbf{y}})}{(12\rho_0 \pi a_0^3)} \ \overline{\left(\frac{\partial p}{\partial t}\right)^2} \, dS(\overline{\mathbf{y}}) \tag{959}$$



Appendix: Curle model

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where $I(\overline{y})$ is the directional acoustic intensity per unit surface.

The model can be enabled for steady and unsteady cases with <u>Reynolds-Averaged Navier-Stokes (RANS)</u> models which can provide turbulence time scale, turbulence length scale and wall shear stress necessary to compute $\overline{(\partial p / (\partial t))}^2$, the mean-square time derivative of the source surface pressure [180].

The acoustic power per unit surface can be reported in dimensional units ($_{W/m^2}$) and in dB:

$$SAP(dB) = 10 \log \frac{SAP}{P_{ref}}$$
(960)

where P_{ref} is the reference acoustic power.



Appendix: Proudman model

The analytical result of Proudman estimated the local acoustic power generated by unit volume of isotropic turbulence having no mean flow.

See [194]. Proudman considered the generation of noise by isotropic turbulence and using statistical models of various two-point moments, using the Lighthill analogy. In Proudman's high-Reynolds model for isotropic turbulence in near incompressible flow, Lilley added the effects of retarded time in th evaluation of the two-point covariance of Lighthill's tensor (an effect previously neglected by Proudman), and obtained the following expression for acoustic power, AP per unit volume:

$$AP = \alpha \rho_0 \frac{u^3}{l} \frac{u^5}{a_0^5}$$
(990)

where α is a constant related to the shape of the longitudinal velocity correlation, u is the root mean square of one of the velocity components, l is the longitudinal integral length scale of the velocity, ρ_0 is the far-field density and a_0 is the far-field sound speed.

In Proudman's original derivation [194], α is approximately 13. Lilley [187] found it to be about 10.96. In a DNS simulation done by Sarkar and Hussain [198] α = 2.6, and in the LES study done by Witkowska and Juve [199] α = 2.5. In Proudman's paper, the terms of u and ε can be written as:

$$u = \sqrt{\left(\frac{2}{3}k\right)} \quad , \qquad \varepsilon = \frac{1.5u^3}{l} \tag{991}$$

In terms of STAR-CCM+ of the turbulence velocity scale and of the turbulence length scale, the local acoustic power due to the unit volume of isotropic turbulence (in W/m^3) becomes:

$$AP = \alpha_c \ \rho_0 \frac{U^3}{L} \frac{U^5}{a_0^5}$$
(992)



Appendix: Proudman model

In terms of STAR-CCM+ of the turbulence velocity scale and of the turbulence length scale, the local acoustic power due to the unit volume of isotropic turbulence (in W/m^3) becomes:

$$AP = \alpha_c \ \rho_0 \frac{U^3}{L} \frac{U^5}{a_0^5}$$
(992)

with:

$$U = \frac{L}{T} \quad , \qquad \alpha_c = 0.629 \tag{993}$$

where ho_0 is the far-field density, U is the turbulence velocity, L is the turbulence length scale, T is the turbulence time scale and a_0 is the far-field sound speed.

The rescaled constant is based on Direct Numerical Simulation for isotropic turbulence done by Sarkar and Hussaini [198].

The total acoustic power per unit volume can be reported in dimensional units ($_{W/m^3}$) and in dB:

$$AP(dB) = 10 \log\left(\frac{AP}{P_{ref}}\right)$$
(994)

where P_{ref} is the reference acoustic power.